Determination of the 232 Th(n, γ) cross section from 10 keV to 200 keV at the **Back-n facility at CSNS**



Jincheng Wang, Jie Ren, Xichao Ruan

Key Laboratory of Nuclear Data, China Institute of Atomic Energy, Beijing, 102413, China



I. Introduction

- Thorium-uranium (Th-U) nuclear fuel cycle, The fissile ²³³U is generated through the ²³²Th neutron capture reaction and the subsequent two consecutive β decays. Compared to the conventional uranium-plutonium cycle, the lighter nucleus ²³²Th in the Th-U cycle is used, significantly reducing the radioactive of the spent fuel by greatly reducing the accumulation of heavy transuranium isotopes^[1,2].
- Large parity-nonconservation effects were observed in the p-wave resonance of 232 Th $^{[3-5]}$.



III. Data Analysis

- ◆ The time-of-flight (TOF) technique.
- The pulse height weighting technique^[10] was applied to determine the number of capture reactions from the measured complex γ -rays cascade spectrum following neutron capture in this work.
- The Saturated Resonance technique (SR) is conveniently applied to calibrate the efficiency of the γ -rays detector and neutron flux monitor.
- The background was determined by a lead sample measurement and detailed Monte Carlo simulations.
- To normalize these background components, the black resonance filter technique was applied.



• Existing experimental data





Decomposition of the spectrum FIG5. measured with the natPb sample into the neutron background and the scattered contribution from in-beam γ rays.

FIG6. ²³²Th spectra with filters in the neutron beam. The three background components have been fitted to match the minimum of the filter dips.

- The correction factor consisted of the normalization correction factor, the threshold correction factor, the dead-time correction factor, and the multiple scattering correction factor.
- In this work, the 232 Th(*n*, γ) average cross section is obtained from the experimental data and the ENDF/B-VIII.0 evaluation data, which can be expressed as

$$\sigma_{Th} = \frac{\langle \sigma_{Th} \rangle}{\langle \sigma_{Au} \rangle} \cdot \sigma_{Au}$$

IV. Summary & Conclusions

• This work also carefully evaluated the uncertainties of PHWT, normalization, background subtraction, correction, and relative measurement. The total uncertainty

Neutron Energy(MeV)

FIG2. Existing experimental data of ²³²Th.

II. Experimental

◆ The back-streaming white neutron beam-line (Back-n) of China Spallation **Neutron Source (CSNS)**

The Back-n of the CSNS became fully operational in March 2018. At the Back-n of the CSNS, neutrons are produced by spallation reactions induced by double bunches per pulsed, 41 ns wide (FWHM) and the interval between the two bunches is 410 ns, 1.6 GeV/c proton beam with a typical repetition rate of 25 Hz, impinging on a tantalumcladded and water-cooled sliced tungsten target.



is in the range of 4.5% - 4.8%.

Table 1 Different components of estimated uncertainty in the measured cross section.

Source of uncertainty	²³² Th (%)	¹⁹⁷ Au (%)
PHWT	1.2	
Normalization	2.0	1.0
Background subtraction	1.5	1.0
Correction	1.3	
Counting statistics	1.0	
Relative measurement	3.2 - 3.9	

- The determined experimental data by taking $^{197}Au(n, \gamma)$ cross sections as reference are consistent with the evaluations of the CENDL-3.2 and JENDL-5 within the uncertainties from 10 to 200 keV.
- ◆ TALYS calculation agrees very well with the experimental data.



detectors was 90 degrees.

FIG4. Photo of the C_6D_6 detector system at Back-n.

10 10² Neutron Energy (keV)

FIG. 7 The capture cross sections of ²³²Th multiplied by the square root of the neutron energy.

V. Acknowledgments

The authors are indebted to the operating crew of the CSNS Back-n white neutron source. Dr. Peter Schillebeeckx from EU-JRC-IRMM is appreciated for the helpful discussions. This work was supported by the National Natural Science Foundation of China (Grant Nos.11790321) and the Youth Talent Program of China National Nuclear Corporation.

[1] W. Gudowski, Nuclear Physics A 752, 623 (2005).

[2] M. Salvatores, I. Slessarev, and V. Berthou, Progress in Nuclear Energy 38, 167 (2001). [3] G. E. Mitchell, J. D. Bowman, and H. A. Weidenmuller, Reviews of Modern Physics **71**, 445 (1999). [4] E. I. Sharapov, J. D. Bowman, B. E. Crawford, P. P. J. Delheij, et al., Phys. Rev. C 61, 025501 (2000). [5] S. L. Stephenson, J. D. Bowman, et al. Phys. Rev. C 58, 1236 (1998).

[6] R. L. Macklin and R. R. Winters, Nuclear Science and Engineering 78, 110 (1981). [7] M. Salvatores, I. Slessarev, and V. Berthou, Progress in Nuclear Energy 38, 167 (2001). [8] A. Borella, K. Volev, A. Brusegan, P. Schillebeeckx, F. Corvi, et al., Nuclear Science and Engineering (2006). [9] G. Aerts, U. Abbondanno, H. 'Alvarez, et al., Phys. Rev. C 73, 054610 (2006). [10] J. Ren, X. Ruan, W. Jiang, J. Bao, J. Wang, Q. Zhang, G. Luan, H. Huang, et al., Chinese Physics C 46, 044002 (2022).

FIG1. The fissile ²³³U is generated through the ²³²Th.